### Input-Output Analysis of Channel Flows

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### Collaborators

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Support:



Energy, Power & Adaptive Systems Program (ECCS)

Control Systems (CMMI)



Dynamics & Control Program

## Context: Networked Systems and DPS



## Context: Networked Systems and DPS



### **Context: Case Studies**



## **Context: Emerging Common Themes**



Network topology imposes asymptotic "hard performance limits"

## Turbulence in Streamlined Flows (Boundary Layers)



boundary layer turbulence



side view



top view

# Turbulence in Streamlined Flows (Boundary Layers)



skin-friction drag: laminar vs. turbulent

- Streamlining a vehicle reduces form drag
- Still stuck with: Skin-Friction Drag (higher in Turbulent BL than in Laminar BL)
- Same in pipe flows (increases required pumping power)

## Control of Boundary Layer Turbulence

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Intuition: must have ability to actuate at spatial scale comparable to flow structures

spatial-bandwidth of controller  $\geq$  plant's bandwidth

CDS20, Aug 2014 8 / 17

## Control of Boundary Layer Turbulence



corrugated skin

compliant skin

Intuition: must have ability to actuate at spatial scale comparable to flow structures

spatial-bandwidth of controller  $\geq$  plant's bandwidth

Caveat: Plant's dynamics are not well understood
 obstacles
 {
 not only device technology
 also: dynamical modeling and control design
 }

#### Mathematical Modeling of Transition: Hydrodynamic Stability

The Navier-Stokes (NS) equations:

$$\partial_t \mathbf{u} = -\nabla_{\mathbf{u}} \mathbf{u} - \operatorname{grad} p + \frac{1}{R} \Delta \mathbf{u}$$
  
 $0 = \operatorname{div} \mathbf{u}$ 



Hydrodynamic Stability:

view NS as a dynamical system

• *laminar flow*  $\bar{\mathbf{u}}_R :=$  a stationary solution of the NS equations (an *equilibrium*)

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view NS as a dynamical system

• *laminar flow*  $\bar{\mathbf{u}}_{R} :=$  a stationary solution of the NS equations (an *equilibrium*)

laminar flow  $\bar{\mathbf{u}}_R$  stable

 $\longleftrightarrow$  i

i.c. 
$$\mathbf{u}(0) \neq \bar{\mathbf{u}}_R$$
,  
 $\mathbf{u}(t) \stackrel{t \to \infty}{\longrightarrow} \bar{\mathbf{u}}_R$ 



- typically done with dynamics linearized about  $\bar{\mathbf{u}}_{R}$
- various methods to track further "non-linear behavior"

#### Mathematical Modeling of Transition: Hydrodynamic Stability

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 $0 = \operatorname{div} \mathbf{u}$ 



Hydrodynamic Stability:

view NS as a dynamical system

• A very successful (*phenomenologically predictive*) approach for many decades

• However: it fails badly in the special (but important) case of streamlined flows

#### Mathematical Modeling of Transition: Adding Signal Uncertainty

Decompose the fields as

Fluctuation dynamics:

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 $\mathbf{u} = \mathbf{\bar{u}}_{R} + \mathbf{\tilde{u}}$   $\uparrow \qquad \uparrow$   $laminar \qquad fluctuations$ In *linear* hydrodynamic stability,  $-\nabla_{\tilde{u}}\tilde{u}$  is ignored

 $\partial_t \tilde{\mathbf{u}} = -\nabla_{\tilde{\mathbf{u}}_R} \tilde{\mathbf{u}} - \nabla_{\tilde{\mathbf{u}}} \tilde{\mathbf{u}}_R - \operatorname{grad} \tilde{p} + \frac{1}{R} \Delta \tilde{\mathbf{u}} - \nabla_{\tilde{u}} \tilde{u} + \mathbf{d}$ 0 = div  $\tilde{\mathbf{u}}$ 

► a time-varying *exogenous disturbance* field d (e.g. random body forces)



Input-Output view of the Linearized NS Equations

Jovanovic, BB, '05 JFM

#### Input-Output Analysis of the Linearized NS Equations

$$\partial_t \begin{bmatrix} \Delta \tilde{v} \\ \tilde{\omega} \end{bmatrix} = \begin{bmatrix} U'' \partial_x - U \Delta \partial_x + \frac{1}{R} \Delta^2 & 0 \\ -U' \partial_z & -U \partial_x + \frac{1}{R} \Delta \end{bmatrix} \begin{bmatrix} \tilde{v} \\ \tilde{\omega} \end{bmatrix} + \begin{bmatrix} -\partial_{xy} & \partial_x^2 + \partial_z^2 & -\partial_{zy} \\ \partial_z & 0 & -\partial_x \end{bmatrix} \begin{bmatrix} d_x \\ d_y \\ d_z \end{bmatrix}$$
$$\begin{bmatrix} \tilde{u} \\ \tilde{v} \\ \tilde{w} \end{bmatrix} = \left( \partial_x^2 + \partial_z^2 \right)^{-1} \begin{bmatrix} \partial_{xy} & -\partial_z \\ \partial_x^2 + \partial_z^2 & 0 \\ \partial_{zy} & \partial_x \end{bmatrix} \begin{bmatrix} \tilde{v} \\ \tilde{\omega} \end{bmatrix}$$



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• eigs  $(\mathcal{A})$ : determine stability

(standard technique in Linear Hydrodynamic Stability)

• Transfer Function  $d \longrightarrow \tilde{u}$ : determines response to disturbances  $\begin{pmatrix} uncommon in Fluid Mechanics \\ an "open system" \end{pmatrix}$ 

#### Input-Output Analysis of the Linearized NS Equations

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Surprises:

Even when A is stable

the gain  $\mathbf{d} \longrightarrow \tilde{\mathbf{u}}$  can be very large ( $(H^2 \text{ norm})^2$  scales with  $R^3$ )

Input-output resonances

very different from least-damped modes of  $\ensuremath{\mathcal{A}}$ 



However: Pole Locations  $\nleftrightarrow$  Frequency Response Peaks Theorem: Given any desired pole locations

 $z_1, \ldots, z_n \in \mathbb{C}_-$  (LHP),

and any stable frequency response  $H(j\omega)$ , arbitrarily close approximation is achievable with

$$\left\| H(s) - \left( \sum_{i=1}^{N_1} \frac{\alpha_{1,i}}{(s-z_1)^i} + \cdots + \sum_{i=1}^{N_n} \frac{\alpha_{n,i}}{(s-z_n)^i} \right) \right\|_{\mathcal{H}^2} \le \epsilon$$

by choosing any of the  $N_k$ 's large enough



However: Pole Locations **Frequency Response Peaks**  $\leftrightarrow$ Theorem: Given any desired pole locations  $z_1, \ldots, z_n \in \mathbb{C}_-$  (LHP), and any stable frequency response  $H(j\omega)$ , arbitrarily close approximation is achievable with  $\left\| H(s) - \left( \sum_{i=1}^{N_1} \frac{\alpha_{1,i}}{(s-z_1)^i} + \dots + \sum_{i=1}^{N_n} \frac{\alpha_{n,i}}{(s-z_n)^i} \right) \right\|_{s=0} \le \epsilon$ by choosing any of the  $N_k$ 's large enough

Remarks:

- No necessary relation between pole locations and system resonances
- (  $\epsilon \to 0 \Rightarrow N_k \to \infty$ ), i.e. this is a *large-scale systems* phenomenon
- Large-scale systems: IO behavior not always predictable from modal behavior

**However:** Pole Locations  $\leftrightarrow$  Frequency Response Peaks MIMO case:  $H(s) = (sI - A)^{-1}$ 

• If A is normal (has orthogonal eigenvectors), then

$$\sigma_{\max}\left(\left(j\omega I - A\right)^{-1}\right) = \frac{1}{_{\text{distance}}\left(j\omega, \text{nearest pole}\right)}$$

 If A is non-normal : no clear relation between singular value plot ↔ eigs(A)





#### Spatio-temporal Impulse and Frequency Responses

Translation invariance in x & z implies

• Impulse Response (Green's Function)

$$\tilde{\mathbf{u}}(t,x,y,z) = \int G(t-\tau,x-\xi,\mathbf{y},\mathbf{y}',z-\zeta) \,\mathbf{d}(\tau,\xi,y',\zeta) \,d\tau d\xi dy' d\zeta$$
$$\tilde{\mathbf{u}}(t,x,..,z) = \int G(t-\tau,x-\xi,z-\zeta) \,\mathbf{d}(\tau,\xi,..,\zeta) \,d\tau d\xi d\zeta$$

$$f(t,x,.,z) = \int \mathcal{G}(t-\tau,x-\xi,z-\zeta) \, \mathbf{d}(\tau,\xi,.,\zeta) \, d\tau d\xi d\zeta$$

 $\mathcal{G}(t, x, z)$  : Operator-valued impulse response

- Frequency Response
  - $\tilde{\mathbf{u}}(\omega, k_x, k_z) = \mathcal{G}(\omega, k_x, k_z) \mathbf{d}(\omega, k_x, k_z)$
  - $\mathcal{G}(\omega, k_x, k_z)$  : Operator-valued frequency response (Packs lots of information!)

• Spectrum of A:

$$\sigma(\mathcal{A}) = \overline{\bigcup_{k_x,k_z} \sigma\left(\hat{\mathcal{A}}(k_x,k_z)\right)}$$



$$\partial_t \Psi = \mathcal{A} \Psi + \mathcal{B} \mathbf{d}$$

$$\tilde{\mathbf{u}} = \mathcal{C} \Psi$$

$$\bullet \text{ IR: } \mathcal{G}(t, x, z)$$

$$\bullet \text{ FR: } \mathcal{G}(\omega, k_x, k_z)$$

**Modal Analysis**: Look for unstable eigs of  $\mathcal{A} \left( \bigcup_{k_x,k_z} \sigma \left( \hat{\mathcal{A}}(k_x,k_z) \right) \right)$ 

Flow type	Classical linear theory R <sub>c</sub>	Experimental R <sub>c</sub>
Channel Flow	5772	$\approx$ 1,000-2,000
Plane Couette	$\infty$	$\approx$ 350
Pipe Flow	$\infty$	pprox 2,200-100,000

$$\partial_{t}\Psi = \mathcal{A}\Psi + \mathcal{B} \mathbf{d}$$

$$\tilde{\mathbf{u}} = \mathcal{C}\Psi$$

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**Modal Analysis**: Look for unstable eigs of  $\mathcal{A} = \left(\bigcup_{k_x,k_z} \sigma\right)$ 





top view

$$\underbrace{\partial_t \Psi = \mathcal{A} \Psi + \mathcal{B} \mathbf{d}}_{(\text{patietergoal system})} \underbrace{\partial_t \Psi = \mathcal{A} \Psi + \mathcal{B} \mathbf{d}}_{\tilde{\mathbf{u}} = \mathcal{C} \Psi}$$
  
• IR:  $\mathcal{G}(t, x, z)$   
• FR:  $\mathcal{G}(\omega, k_x, k_z)$ 

**Modal Analysis**: Look for unstable eigs of A (

$$\left(\bigcup_{k_x,k_z}\sigma\left(\hat{\mathcal{A}}(k_x,k_z)\right)\right)$$











Impulse Response Analysis: Channel Flow @ R = 2000





#### similar to "turbulent spots"

Jovanovic, BB, '01 ACC,

more movies and pics at http://engineering.ucsb.edu/~bamieh/pics/impulse\_page.html

 $\mathcal{G}(\omega, k_x, k_z)$  is a *LARGE* object! (very "data rich"!)



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one visualization method:  $\sup_{\omega} \sigma_{\max} \left( \mathcal{G}(\omega, k_x, k_z) \right)$ 



What do the corresponding flow structures look like?



streamwise velocity isosurfaces

streamwise vorticity isosurfaces

 $\mathcal{G}(\omega, k_x, k_z)$  is a *LARGE* object! (very "data rich"!)

one visualization method:  $\sup_{\omega} \sigma_{\max} \left( \mathcal{G}(\omega, k_x, k_z) \right)$ 



What do the corresponding flow structures look like?

much closer (than TS waves) to structures seen in turbulent boundary layers







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